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## Letter

# BAlN alloy for enhanced two-dimensional electron gas characteristics of GaN/AlGaN heterostructures

Rongyu Lin<sup>1</sup> , Xinwei Liu<sup>1</sup>, Kaikai Liu<sup>1</sup> , Yi Lu<sup>1</sup>, Xinke Liu<sup>2</sup>  and Xiaohang Li<sup>1</sup> 

<sup>1</sup> King Abdullah University of Science and Technology (KAUST), Advanced Semiconductor Laboratory, Thuwal 23955-6900, Saudi Arabia

<sup>2</sup> College of Materials Science and Engineering, College of Electronics and Information Engineering, Guangdong Research Center for Interfacial Engineering of Functional Materials, Shenzhen Key Laboratory of Special Functional Materials, Chinese Engineering and Research Institute of Microelectronics, Shenzhen University, Shenzhen 518060, People's Republic of China

E-mail: [xkliu@szu.edu.cn](mailto:xkliu@szu.edu.cn) and [xiaohang.li@kaust.edu.sa](mailto:xiaohang.li@kaust.edu.sa)

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## Abstract

Emerging wide bandgap BAlN alloys have potential for improved III-nitride power devices, including high electron mobility transistors (HEMTs). Yet, few relevant studies have been carried out. In this work, we have investigated the use of the  $B_{0.14}Al_{0.86}N$  alloy as part or the entirety of the interlayer between the GaN buffer and the AlGaN barrier in the conventional GaN/AlGaN heterostructure. The numerical results show considerable improvement of the two-dimensional electron gas (2DEG) concentration with small 2DEG leakage into the ternary layer by replacing the conventional AlN interlayer by either the  $B_{0.14}Al_{0.86}N$  interlayer or the  $B_{0.14}Al_{0.86}N$ /AlN hybrid interlayer. Consequently, the transfer characteristics can be improved. The saturation current can be enhanced as well. For instance, the saturation currents for HEMTs with the 0.5 nm  $B_{0.14}Al_{0.86}N$ /0.5 nm AlN hybrid interlayer and the 1 nm  $B_{0.14}Al_{0.86}N$  interlayer are 5.8% and 2.2% higher than that for the AlN interlayer when  $V_{GS} - V_{th} = +3$  V.

**Keywords:** GaN, high electron mobility transistor, BAlN, interlayer, two-dimensional electron gas

(Some figures may appear in colour only in the online journal)

## 1. Introduction

In the past decades, III-nitride semiconductors, especially AlGaIn, InGaIn, and GaN, have been widely employed for optical and electronic devices due to their superior properties,

such as suitable bandgap, high saturation velocity, high breakdown field, and high chemical and thermal stability. Invented in 1993 based on the two-dimensional electron gas (2DEG) at the AlGaIn/GaN heterointerface induced by the net polarization charge, the GaN-based high electron mobility transistor (HEMT) has attracted enormous research interest due to its vast high power and high speed applications [1]. Since then, substantial efforts have been made to improve the GaN-based HEMT performance [2]. Previously, there have been reports of the use of organic interlayers to regulate the



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charge transport in organic transistors [3–6], and the organic molecules could also be applied in GaN-based HEMTs to improve the transport properties [7]. In general, a higher Al content of the AlGa<sub>N</sub> barrier layer can induce larger net polarization charge at the AlGa<sub>N</sub>/GaN heterointerface, resulting in higher 2DEG concentration in the conventional AlGa<sub>N</sub>/GaN HEMT structure. However, the electron mobility can be compromised due to the scattering caused by the alloy disorder. Additionally, further increasing the Al content of the AlGa<sub>N</sub> barrier layer could degrade the interface quality due to the larger lattice mismatch between the AlGa<sub>N</sub> barrier layer and the GaN buffer layer [8].

To overcome the issues, the insertion of a thin ( $\leq 2$  nm) AlN interlayer between the barrier layer and the GaN channel layer has been proposed and implemented, partly because the binary AlN does not cause alloy scattering. A thicker AlN interlayer is not feasible, as the AlN interlayer could relax due to the large lattice mismatch with the GaN buffer layer. Furthermore, the AlN interlayer can elevate the barrier height for the 2DEG due to its higher conduction band edge than that of the AlGa<sub>N</sub> barrier. Since the penetration of the electron wave function into a given barrier is in inverse proportion to the barrier height, the AlN interlayer could suppress the 2DEG leakage into the barrier, further suppressing the alloy scattering [9]. Moreover, the net polarization charge at the AlN/GaN heterointerface is greater than that of the AlGa<sub>N</sub>/GaN heterointerface, which induces a higher 2DEG concentration. Combining the three benefits, the AlN interlayer has been employed to minimize the alloy scattering and enhance the 2DEG concentration [10].

The wurtzite BAlN alloys with large bandgap have application potentials for III-nitride devices. Recently, meaningful progress has been made in studying epitaxy and the properties of BAlN thin films and heterojunctions. The growth of a highly reflective B<sub>x</sub>Al<sub>1-x</sub>N/AlN distributed Bragg reflector was conducted by metalorganic vapor phase epitaxy (MOVPE) [11]. The growth of five-period B<sub>x</sub>Al<sub>1-x</sub>N/AlN heterostructures with the boron content of 11% has been demonstrated [12], and researchers have performed the growth of single-crystalline BAlN layers with relatively high boron composition by MOVPE [13–15]. Moreover, the lattice constant and polarization properties of the BAlN alloys have been studied, in which large variations of those important quantities were found [16–18]. The net heterointerface polarization charges of the strained Al<sub>x</sub>Ga<sub>1-x</sub>N/GaN ( $0 \leq x \leq 1$ ) heterostructures on relaxed GaN are much smaller compared to those of the strained Al<sub>x</sub>Ga<sub>1-x</sub>N/B<sub>0.14</sub>Al<sub>0.86</sub>N ( $0 \leq x \leq 1$ ) heterostructures on relaxed GaN with similar Al composition  $x$ 's of Al<sub>x</sub>Ga<sub>1-x</sub>N. For instance, the absolute values of the net heterointerface polarization charges of the Al<sub>0.3</sub>Ga<sub>0.7</sub>N/GaN heterojunction is  $0.0335 \text{ C m}^{-2}$  [19]. In comparison, the absolute values of the net heterointerface polarization charges of the Al<sub>0.3</sub>Ga<sub>0.7</sub>N/B<sub>0.14</sub>Al<sub>0.86</sub>N and GaN/B<sub>0.14</sub>Al<sub>0.86</sub>N heterojunction are  $0.1364 \text{ C m}^{-2}$  and  $0.1699 \text{ C m}^{-2}$ , which could offer significantly more band bending. Furthermore, the band alignments of the BAlN/(Al)Ga<sub>N</sub> heterojunctions have been determined, showing an extremely large conduction band offset between BAlN and (Al)Ga<sub>N</sub> due to the type-II nature, which is distinct from the type-I nature for the Al<sub>x</sub>Ga<sub>1-x</sub>N/GaN

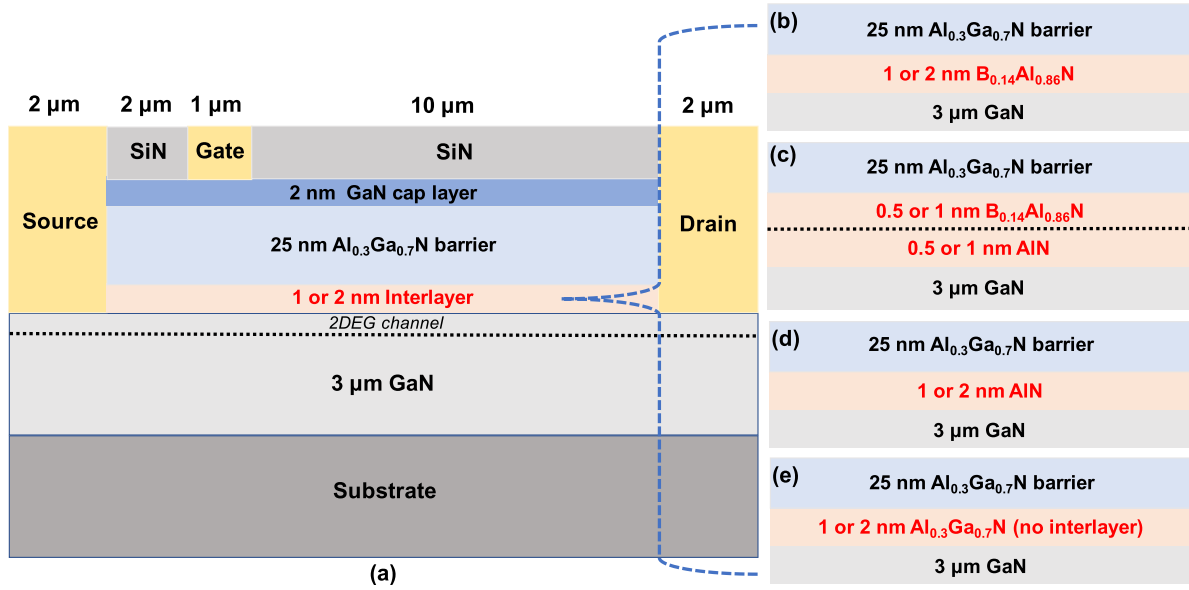
( $0 \leq x \leq 1$ ) heterostructures [20, 21]. Despite the unique polarization and band offset properties, however, there have been few device research reports about the incorporation of the BAlN alloys in the technically-important HEMTs.

In this work, we propose to incorporate the BAlN alloy into the GaN HEMTs by leveraging its unique band alignment and polarization properties. Specifically, a thin BAlN layer is employed as the whole interlayer or as part of the interlayer in the conventional AlGa<sub>N</sub>/GaN-based HEMTs to provide considerably larger polarization charge at the heterointerface and thus 2DEG concentrations, while effectively suppressing the electron leakage into the AlGa<sub>N</sub> barrier layer through a larger conduction band offset than the conventional AlN interlayer.

## 2. Impact of interlayers on 2DEG characteristics

The simulation structures are constructed based on the conventional wurtzite III-polar GaN-based HEMT comprising a fully-relaxed GaN buffer layer, a 1 or 2 nm interlayer, and an Al<sub>0.3</sub>Ga<sub>0.7</sub>N barrier layer, shown in figure 1(a) [22–24]. For the interlayers, various designs under consideration include: 1 nm B<sub>0.14</sub>Al<sub>0.86</sub>N interlayer, 2 nm B<sub>0.14</sub>Al<sub>0.86</sub>N interlayer, 0.5 nm B<sub>0.14</sub>Al<sub>0.86</sub>N/0.5 nm AlN hybrid interlayer, 1 nm B<sub>0.14</sub>Al<sub>0.86</sub>N/1 nm AlN hybrid interlayer, 1 nm AlN interlayer, and 2 nm AlN interlayer, as shown in figures 1(b)–(e). It should be noted that for the structures without the interlayer, we increase the Al<sub>0.3</sub>Ga<sub>0.7</sub>N barrier thickness from 25 nm to 26 nm and 27 nm corresponding to the 1 and 2 nm Al<sub>0.3</sub>Ga<sub>0.7</sub>N interlayers, respectively, in order to keep the total thickness fixed for fair comparison. The B-content of 14% is chosen, i.e. the B<sub>0.14</sub>Al<sub>0.86</sub>N alloy, because its band alignment with (Al)Ga<sub>N</sub> could be promising for the enhanced GaN-based HEMT and it has been realized epitaxially by our group. As shown by previous studies, sharp interfaces were observed for B<sub>0.14</sub>Al<sub>0.86</sub>N/Al<sub>0.7</sub>Ga<sub>0.3</sub>N and B<sub>0.14</sub>Al<sub>0.86</sub>N/GaN heterojunctions without boron diffusion, which have much larger thickness than our proposed 1 nm and 0.5 nm of the B<sub>0.14</sub>Al<sub>0.86</sub>N layers of the hybrid B<sub>0.14</sub>Al<sub>0.86</sub>N/AlN interlayers [20, 21]. The maximum thickness of the interlayer is limited at 2 nm as larger thicknesses could lead to lattice relaxation given the large lattice mismatch between the interlayer and the buffer layer, thus poorer material quality [25]. The entire HEMT structure is assumed to be fully strained to the GaN buffer layer. The B<sub>0.14</sub>Al<sub>0.86</sub>N/AlN heterojunction where AlN situates between B<sub>0.14</sub>Al<sub>0.86</sub>N and GaN is employed to minimize potential alloy scattering. The dielectric SiN<sub>x</sub> is applied for surface passivation. The length of the Schottky gate, the gate-to-source spacing, and the gate-to-drain spacing are set according to the typical values of the GaN-based HEMT shown in figure 1(a). The channel dimensions are the same for all the structures for fair comparison.

The simulation employs APSYS developed by Crosslight Inc. [26]. The electrical properties of the structures are performed by solving Poisson's equation and the continuity equation. The transport model of electrons and holes including drift and diffusion are considered. The electron transport is based on the mobility models inside APSYS, which are the



**Figure 1.** (a) Cross-sectional schematics of the HEMT structures with different interlayer configurations: (b) 1 or 2 nm B<sub>0.14</sub>Al<sub>0.86</sub>N, (c) 0.5 nm B<sub>0.14</sub>Al<sub>0.86</sub>N/0.5 nm AlN or 1 nm B<sub>0.14</sub>Al<sub>0.86</sub>N/1 nm AlN, (d) 1 or 2 nm AlN, (e) 1 or 2 nm Al<sub>0.3</sub>Ga<sub>0.7</sub>N (i.e. no interlayer).

Caughey–Thomas model and modified transferred electron model for different electric field conditions [27, 28]. The formation of the 2DEG amid the HEMT structures are mainly related to the heterointerface polarization difference and the band offset. The lattice constants, polarization parameters, and band offsets of the involved materials are from [16, 17, 20, 21], from which it can be deduced that the B<sub>0.14</sub>Al<sub>0.86</sub>N/GaN heterojunction exhibits a significantly larger conduction band offset and heterointerface polarization difference as compared to the conventional AlN/GaN and Al<sub>0.3</sub>Ga<sub>0.7</sub>N/GaN heterojunctions, potentially leading to much higher 2DEG concentration and less 2DEG leakage from the channel into the barrier.

The conduction band diagrams and the polarization-induced electron concentration of the HEMT structures with and without the 1 nm and 2 nm interlayers under zero bias are plotted in figures 2(a) and (d). Due to larger polarization sheet charge induced by B<sub>0.14</sub>Al<sub>0.86</sub>N, the conduction bands of the GaN buffer layer bend down the most below the Fermi level at the interlayer/buffer interface for the 1 nm and 2 nm B<sub>0.14</sub>Al<sub>0.86</sub>N interlayers from figures 2(a) and (c), respectively, followed by the hybrid interlayers and the AlN interlayers. The structures without the interlayers have the smallest conduction band bending. As the increased band bending below the Fermi level induces more electrons at the channel, the use of the B<sub>0.14</sub>Al<sub>0.86</sub>N interlayers leads to the highest concentration of electrons, followed by the hybrid interlayers, the AlN interlayers, and the structures without the interlayers are shown in figures 2(b) and (d).

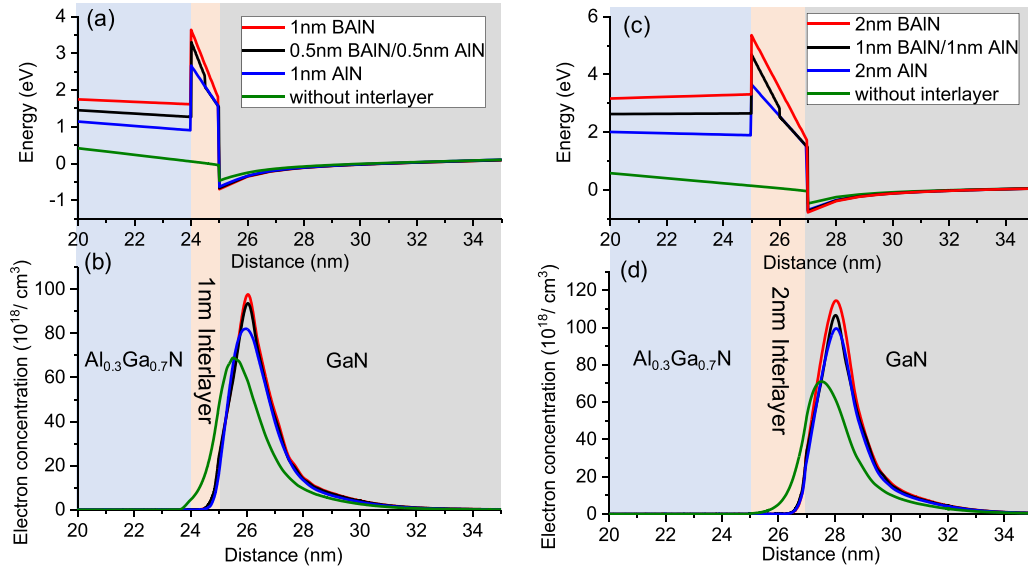
The corresponding 2DEG sheet densities of different structures are shown by red squares in figure 3. The structures with the 1 nm B<sub>0.14</sub>Al<sub>0.86</sub>N interlayer and the 1 nm hybrid interlayer have 2DEG sheet densities of  $1.834 \times 10^{13}$  and  $1.710 \times 10^{13} \text{ cm}^{-2}$ , which are 11.6% and 4.0% higher than  $1.644 \times 10^{13} \text{ cm}^{-2}$  for the 1 nm AlN interlayer, respectively. Additionally, the 2DEG sheet densities of the structures with

the 2 nm B<sub>0.14</sub>Al<sub>0.86</sub>N and hybrid interlayers are  $2.120 \times 10^{13}$  and  $2.001 \times 10^{13} \text{ cm}^{-2}$ , which are 13.6% and 7.2% larger than  $1.866 \times 10^{13} \text{ cm}^{-2}$  for the 2 nm AlN interlayer. The structures without the interlayers, i.e. 1 nm Al<sub>0.3</sub>Ga<sub>0.7</sub>N and 2 nm Al<sub>0.3</sub>Ga<sub>0.7</sub>N in figure 3, have the lowest 2DEG sheet densities of  $1.532 \times 10^{13}$  and  $1.517 \times 10^{13} \text{ cm}^{-2}$ , respectively.

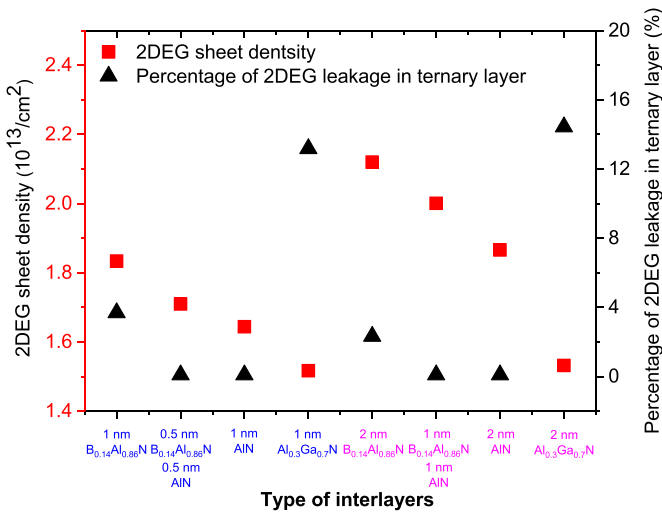
Though the B<sub>0.14</sub>Al<sub>0.86</sub>N interlayer has shown a considerably higher 2DEG sheet density, the leakage of the 2DEG into the B<sub>0.14</sub>Al<sub>0.86</sub>N ternary alloy could lead to scattering induced by the alloy disorder, which could compromise the 2DEG mobility and the device characteristics [29]. Meanwhile, figures 2(a) and (c) show that the structures with the 1 nm and 2 nm B<sub>0.14</sub>Al<sub>0.86</sub>N interlayers exhibit the highest conduction band edge in the respective figures contributing to large effective conduction band offsets of 4.33 eV and 6.13 eV between the interlayer and the channel, respectively, which is desirable for the suppressed 2DEG leakage. The B<sub>0.14</sub>Al<sub>0.86</sub>N/AlN hybrid interlayers result in a slightly lower effective conduction band offset but still higher than that by the conventional AlN interlayers in figures 2(a) and (c). The structures without the interlayer have drastically lower conduction band edge and hence the smallest effective conduction band offset.

Figures 2(b) and (d) manifest the spatial distribution of the 2DEG in the vicinity of the GaN channel. The majority of the 2DEG of all the structures are situated amid the GaN channel, though the structures without the interlayers have the largest leakage of the 2DEG in the Al<sub>0.3</sub>Ga<sub>0.7</sub>N ternary layer. To quantify the leakage into the ternary layer, i.e. Al<sub>0.3</sub>Ga<sub>0.7</sub>N or B<sub>0.14</sub>Al<sub>0.86</sub>N, the percentage of the 2DEG leakage to the ternary layer is defined as follows and the corresponding values are shown by the black triangles in figure 3.

$$\text{Percentage of 2DEG leakage} = \frac{\text{2DEG in ternary layer}}{\text{Total 2DEG}}. \quad (1)$$



**Figure 2.** (a) Band diagram and (b) electron concentration of the HEMT structures with and without the 1 nm interlayers. (c) Band diagram and (d) electron concentration of the HEMT structures with and without the 2 nm interlayers.



**Figure 3.** The 2DEG sheet density (red) and the percentage of 2DEG leakage in the ternary layer (black) of the HEMT structures with different configurations of 1 nm interlayers (blue axis) and 2 nm interlayers (purple axis).

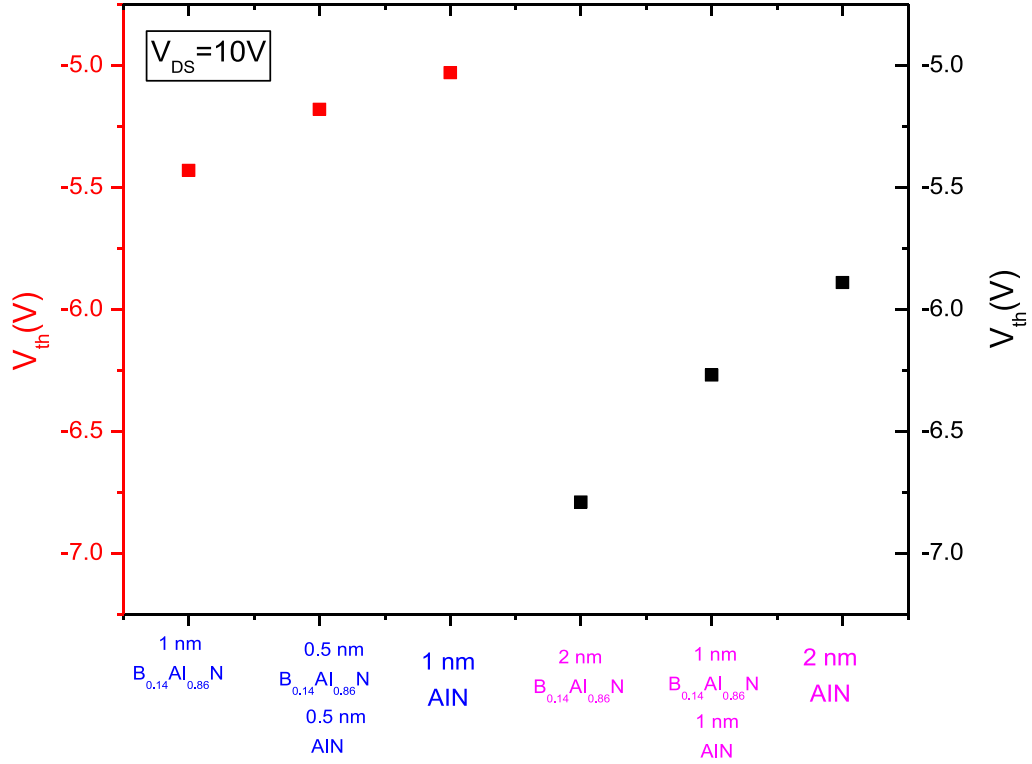
Large percentages of 13.2% and 14.4% are found for the 1 nm and 2 nm  $\text{Al}_{0.3}\text{Ga}_{0.7}\text{N}$  (i.e. no interlayer), which can lead to severe alloy scattering. The percentages for the 1 nm and 2 nm AlN interlayers are less than 0.1%, which show effective leakage suppression. Small percentages less than 0.1% are also found for the 1 nm and 2 nm  $\text{B}_{0.14}\text{Al}_{0.86}\text{N}$ /AlN hybrid interlayers, indicating excellent suppression with larger 2DEG sheet densities as compared with the AlN interlayers. For the 1 nm and 2 nm  $\text{B}_{0.14}\text{Al}_{0.86}\text{N}$  interlayers with the largest 2DEG sheet densities, relatively small percentages of 3.7% and 2.3% are found.

### 3. HEMT device characteristics

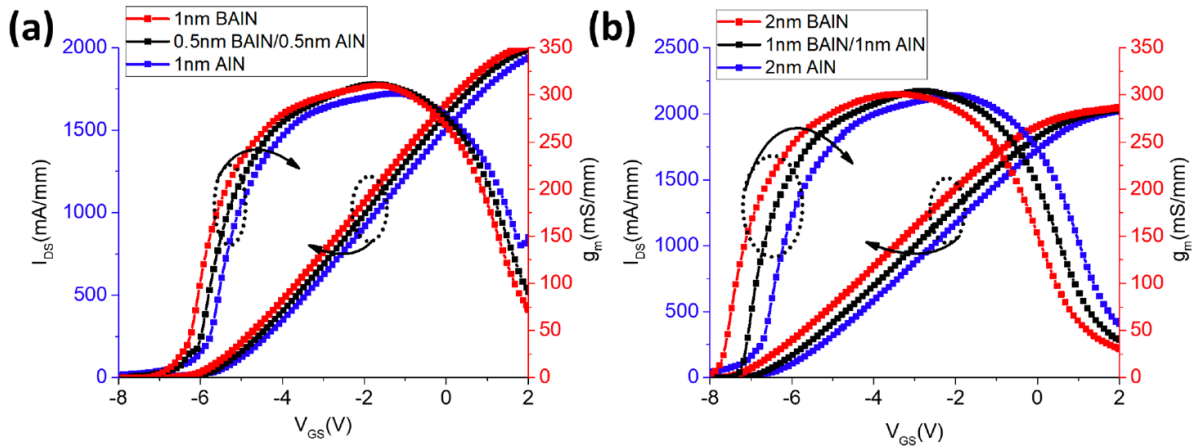
The HEMT device characteristics are investigated subsequent to the study of the 2DEG characteristics. From figure 4, the threshold voltages under  $V_{\text{DS}} = 10$  V for the HEMTs with the interlayers of 1 nm  $\text{B}_{0.14}\text{Al}_{0.86}\text{N}$ , 0.5 nm  $\text{B}_{0.14}\text{Al}_{0.86}\text{N}$ /0.5 nm AlN, and 1 nm AlN are  $-5.43$  V,  $-5.18$  V, and  $-5.03$  V, respectively, which are inversely proportional to the 2DEG sheet density in figure 4. Similarly, for the 2 nm interlayer designs, the threshold voltages are  $-6.79$  V,  $-6.27$  V, and  $-5.89$  V for the 2 nm  $\text{B}_{0.14}\text{Al}_{0.86}\text{N}$ , 1 nm  $\text{B}_{0.14}\text{Al}_{0.86}\text{N}$ /1 nm AlN, and 2 nm AlN interlayers. The overall decreased threshold voltages for the 2 nm interlayer designs are attributed to the higher 2DEG sheet densities, in comparison with the 1 nm interlayer designs shown in figure 4.

The transfer characteristics under  $V_{\text{DS}} = 10$  V are shown in figure 5. The 0.5 nm  $\text{B}_{0.14}\text{Al}_{0.86}\text{N}$ /0.5 nm AlN hybrid interlayer exhibits the maximum transconductance  $g_m$  of  $311.36 \text{ ms mm}^{-1}$ , compared with  $310.00 \text{ ms mm}^{-1}$  for the 1 nm  $\text{B}_{0.14}\text{Al}_{0.86}\text{N}$  interlayer and  $301.6 \text{ ms mm}^{-1}$  for the 1 nm AlN interlayer. Since the gate length remains the same for all structures of interest, the highest maximum transconductance for the 0.5 nm  $\text{B}_{0.14}\text{Al}_{0.86}\text{N}$ /0.5 nm AlN hybrid interlayer indicates improved carrier transport attributed to the minimized 2DEG leakage into the ternary layer and the higher 2DEG sheet density shown in figures 2 and 3. Despite with the 2DEG leakage of 3.7%, the 1 nm  $\text{B}_{0.14}\text{Al}_{0.86}\text{N}$  interlayer design yields a larger maximum transconductance as opposed to the conventional 1 nm AlN interlayer. Also the transconductance of  $310.00 \text{ ms mm}^{-1}$  for the 1 nm  $\text{B}_{0.14}\text{Al}_{0.86}\text{N}$  interlayer is barely lower than that for the 1 nm hybrid interlayer, which could be attributed to the larger 2DEG sheet density. For the structures with the 2 nm interlayers, the hybrid interlayer comprising 1 nm  $\text{B}_{0.14}\text{Al}_{0.86}\text{N}$ /1 nm AlN shows the largest maximum





**Figure 4.** The threshold voltage  $V_{th}$  of the HEMT structures with different interlayers under  $V_{DS} = 10$  V.

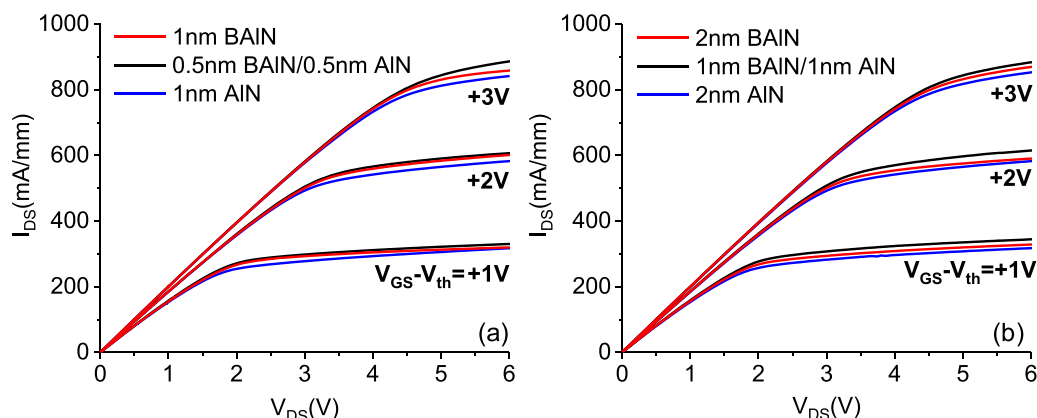


**Figure 5.** Transfer characteristics under  $V_{DS} = 10$  V of the HEMT structures comprising interlayers of (a) 1 nm  $B_{0.14}Al_{0.86}N$ , 0.5 nm  $B_{0.14}Al_{0.86}N/0.5$  nm AlN, and 1 nm AlN; and (b) 2 nm  $B_{0.14}Al_{0.86}N$ , 1 nm  $B_{0.14}Al_{0.86}N/1$  nm AlN, and 2 nm AlN.

transconductance of  $304.13 \text{ ms mm}^{-1}$  among the three, compared to the 2 nm  $B_{0.14}Al_{0.86}N$  interlayer ( $300.50 \text{ ms mm}^{-1}$ ) and the 2 nm AlN interlayer ( $299.63 \text{ ms mm}^{-1}$ ). Again, the  $B_{0.14}Al_{0.86}N$ -containing interlayers lead to higher transconductance than the AlN interlayer for the 2 nm interlayer design. The larger transconductances show the feasibility of employing  $B_{0.14}Al_{0.86}N$  as the entirety or part of the interlayer to enhance the channel carrier transport.

Figures 6(a) and (b) illustrate the  $I_{DS}-V_{DS}$  characteristics of the HEMT structures with different 1 and 2 nm interlayers with varying  $V_{GS}-V_{th}$  of +1, +2, and +3 V. In all cases of figures 6(a) and (b), the hybrid  $B_{0.14}Al_{0.86}N/AlN$  interlayer results in higher saturation current under the same  $V_{GS}-V_{th}$ , with

the HEMTs with the  $B_{0.14}Al_{0.86}N$  interlayer and the AlN interlayer being the second and the third, respectively. For instance, the saturation currents for the  $B_{0.14}Al_{0.86}N/AlN$  hybrid interlayer and the  $B_{0.14}Al_{0.86}N$  interlayer are 5.8% and 2.2% higher than for the AlN interlayer when  $V_{GS}-V_{th} = +3$  V in figure 6(a) for the 1 nm interlayers. For the 2 nm interlayers in figure 6(b), 3.7% and 1.9% saturation current enhancements are observed in the hybrid  $B_{0.14}Al_{0.86}N/AlN$  interlayer and the  $B_{0.14}Al_{0.86}N$  interlayer compared with that for the AlN interlayer under the same  $V_{GS}-V_{th}$ . The enhancements are consistent with the enhanced maximum transconductance in figures 5(a) and (b) for the HEMTs comprising the  $B_{0.14}Al_{0.86}N$ -based interlayers.



**Figure 6.** The  $I_{DS}$ – $V_{DS}$  characteristics of the HEMT structures comprising interlayers of (a) 1 nm  $B_{0.14}Al_{0.86}N$ , 0.5 nm  $B_{0.14}Al_{0.86}N/0.5$  nm AlN, and 1 nm AlN, and (b) 2 nm  $B_{0.14}Al_{0.86}N$ , 1 nm  $B_{0.14}Al_{0.86}N/1$  nm AlN, and 2 nm AlN.

## 4. Conclusion

We have proposed the application of the  $B_{0.14}Al_{0.86}N$  alloy for the interlayers of GaN-based HEMT device structures. The proposed structures comprise either  $B_{0.14}Al_{0.86}N$  interlayers or  $B_{0.14}Al_{0.86}N/AlN$  hybrid interlayers of 1 nm or 2 nm thick. The proposed structures lead to considerable enhancement of the 2DEG sheet density as opposed to the conventional AlN interlayers and result in good 2DEG confinement with low 2DEG leakage into the ternary layer. For instance, the structures with the 1 nm  $B_{0.14}Al_{0.86}N$  interlayer and the 1 nm hybrid interlayer have 2DEG sheet densities of  $1.834 \times 10^{13}$  and  $1.710 \times 10^{13} \text{ cm}^{-2}$ , which are 11.6% and 4.0% higher than  $1.644 \times 10^{13} \text{ cm}^{-2}$  for the 1 nm AlN interlayer, respectively. Additionally, the small percentages of the 2DEG leakage of less than 0.1% are found for the 1 and 2 nm  $B_{0.14}Al_{0.86}N/AlN$  hybrid interlayers, indicating excellent suppression with larger 2DEG sheet densities. The enhanced 2DEG characteristics thanks to the use of the  $B_{0.14}Al_{0.86}N$  alloy lead to higher maximum transconductance and saturation currents than the HEMTs with the conventional AlN interlayers. As this might be the first work applying BAIN into the III-nitride HEMT, we believe that there is still huge potential to optimize the HEMT structure with different boron compositions, or using BAIN as the barrier or back barrier layer. Moreover, the experiences of incorporating BAIN alloys into the HEMTs could be transferred to other optoelectronic devices due to their unique polarization and band structure properties, such as UV LEDs, lasers, and photodetectors.

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## ORCID iDs

Rongyu Lin  <https://orcid.org/0000-0001-5780-2838>  
 Kaikai Liu  <https://orcid.org/0000-0002-1360-9646>  
 Xinke Liu  <https://orcid.org/0000-0002-7648-3972>  
 Xiaohang Li  <https://orcid.org/0000-0002-4434-365X>

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