Full Length Article

Band alignment of orthorhombic Ga$_2$O$_3$ with GaN and AlN semiconductors

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Abstract

G$_2$O$_3$ semiconductors have attracted tremendous research interests because of their fascinating material properties for future-generation energy, electronic, and optoelectronic applications. In the present study, we have performed the epitaxial growth of tin-doped G$_2$O$_3$ on sapphire, GaN, and AlN templates by the pulsed laser deposition technique. The initial characterizations show a two-dimensional mode of single-crystalline orthorhombic G$_2$O$_3$ (κ-G$_2$O$_3$) growth on these substrates with smooth surface morphology. Integrating κ-G$_2$O$_3$ with nitride semiconductors is interesting since both these materials possess polarization, which could induce 2-dimensional carrier gas (2DCG) at the interface. X-ray photoelectron spectroscopy studies reveal that both κ-G$_2$O$_3$/GaN and κ-G$_2$O$_3$/AlN heterostructure form a type-I band structure where the conduction band offset (CBO) was calculated to be 1.38 eV and 1.04 eV, respectively. This unique band alignment with high CBO could lead to the development of efficient power devices.

1. Introduction

Power device technology plays a central role in numerous crucial applications, including electric vehicles, solar farms, 5G/6G, and data centers [1–3]. GaN and AlN are promising wide bandgap (WBG) semiconductors that lead to power devices which operate with considerably higher voltage and efficiency than conventional silicon with smaller bandgap and breakdown field [4]. With decades of development, nitride-based power devices have entered wide market adoption. However, despite that the G$_2$O$_3$ ultra-wide bandgap (UWBG) semiconductor possesses an even greater Baliga figure of merit [5], the research and commercialization of G$_2$O$_3$-based power devices are still at the relatively early stage due to some serious issues such as a lack of effective p-doping techniques and poor thermal conductivity [4].

There are mainly different polymorphs for G$_2$O$_3$, comprising corundum (α), monoclinic (β), defective spinel (γ), cubic (δ), hexagonal (ε), and orthorhombic (κ) phase [6]. One of the polymorphs of G$_2$O$_3$, β-G$_2$O$_3$, has been well explored concerning the growth, material properties, and devices because of its high thermal stability [7–9]. However, the other polymorph of G$_2$O$_3$ are still behind as most of them are in a metastable state and change to a stable β phase at higher temperatures (>500 °C)[10]. Interestingly, in the orthorhombic phase, one of the metastable phases of G$_2$O$_3$, represented as κ-G$_2$O$_3$, shows spontaneous polarization [11,12] and high-temperature stability till 1000°C[13]. Lately, κ-G$_2$O$_3$ has captured huge research attention due to the predicted high value of its spontaneous polarization i.e., 0.23 C/m$^2$ [14,15]. Therefore, this promising property can effectively control the two-dimensional carrier gas (2DCG) at the interfaces between κ-G$_2$O$_3$ and other suitable semiconductors without any chemical impurity doping. Hence, integrating high polarization semiconductors such as nitride (spontaneous polarization of AlN $\approx$ 1.33 C/m; GaN $\approx$ 1.33 C/m) [16] with κ-G$_2$O$_3$ semiconductors can lead to huge polarization difference which opens a prospect to control the interfacial conductivity similar to an AlGaN/GaN heterostructure [17,18]. Furthermore, as compared to nitride semiconductors, the κ-G$_2$O$_3$ phase could show not only the polarization but also the ferroelectric property. This suggests that the spontaneous polarization of κ-G$_2$O$_3$ could be changed by applying an external electric field [19].

Therefore, integrating existing nitride technology with κ-G$_2$O$_3$ semiconductors can provide a favorable platform for exploring emerging phenomena, which motivates new concepts for fundamental physics and applied research. Moreover, this combination can effectively alleviate the thermal issue of κ-G$_2$O$_3$. Thereby, fundamental properties of κ-G$_2$O$_3$ based heterostructures, such as the energy band diagram, band alignment, built-in interface potential, and polarity discontinuity, are crucial for the tuning of polar oxide interfaces toward advanced devices [20]. We have previously studied the band alignment of β-G$_2$O$_3$ and AlN, and a type-II heterostructure has been witnessed [21]. In this work, the study of valence and conduction band offsets (VBO and CBO) of κ-G$_2$O$_3$/GaN and κ-G$_2$O$_3$/AlN heterojunctions have been calculated.
confirmed. The band offset of experimentally and type-I band alignments for both combinations are extremely important because it ascertains the energy barriers in electron and hole transport for the efficient working of electronic and optical devices.

2. Experimental section

In this work, we perform the pulsed laser deposition (PLD) technique for the growth of \( \kappa \)-Ga\(_2\)O\(_3\) on sapphire, GaN, and AlN templates. The 3 \( \mu \)m epitaxial GaN and AlN growths were carried out in a Taiyo Nippon Sanso SR-4000HT horizontal Metal Organic Chemical Vapor Deposition (MOCVD) system on a c-plane sapphire substrate. Trimethylaluminum (TMA), trimethylgallium (TMG), and ammonia (NH\(_3\)) were used as precursors with nitrogen (N\(_2\)) as the carrier gas. The substrate growth temperature is 1340°C and 1150°C while V/III ratio is 742 and 2423 for AlN and GaN growth respectively. Prior to the PLD deposition, the sapphire substrate, MOCVD grown GaN and AlN templates were cleaned ultrasonically with acetone (5 min) followed by isopropyl alcohol (5 min) and then rinsed with deionized water. The thin doped (1.5 wt%) Ga\(_2\)O\(_3\) target was irradiated using a KrF excimer laser (\( \lambda = 248 \) nm) at 100 mJ (on the target) for 15,000 pulses at a frequency of 2 Hz for bulk samples with a thickness of 360 nm. Further, 3 nm thick \( \kappa \)-Ga\(_2\)O\(_3\) was grown on GaN and AlN template for high-resolution X-ray photoemission spectroscopy (HR-XPS) heterointerface analysis. The growth was performed at a substrate temperature of 650°C without the presence of oxygen partial pressure (chamber vacuum: 6x10\(^{-7}\) Torr). The crystal quality of the PLD grown \( \kappa \)-Ga\(_2\)O\(_3\) on the various substrates has been analyzed by X-Ray diffractometer (XRD) and high-resolution scanning transmission electron microscopy (HR-STEM).

3. Results and discussion

Fig. 1 shows the XRD 2theta-omega scan profiles of \( \kappa \)-Ga\(_2\)O\(_3\) layers directly grown on (a) c-plane Sapphire (0001), (b) AlN (0001), and (c) GaN (0001) substrates. \( \kappa \)-Ga\(_2\)O\(_3\) films were grown along (001) direction for all the substrates where the 2theta angle positioned at 19.15°, 38.80°, and 59.75° represents the (002), (004) and (006) plane of diffraction of \( \kappa \)-Ga\(_2\)O\(_3\). The 2theta scan is also represented in logarithmic scale to see each reflection in detail (supplementary information S1). The sharp and intense X-Ray diffraction peaks of \( \kappa \)-Ga\(_2\)O\(_3\) divulge the single crystalline hetero-epitaxial growth on sapphire, GaN, and AlN templates. Further, 2theta positions perceived at 34.5° and 36° belong to GaN and AlN, respectively. All the XRD peak positions are calibrated with a sapphire substrate peak at 41.6°. To confirm the orthorhombic phase of Ga\(_2\)O\(_3\), \( \phi \) scan for a 20 value of 33.05° and a y value of 54.67° is performed (supplementary information S1). At this particular angle, only (122) reflex of the orthorhombic lattice can be observed [22,23]. The 12 peaks correspond to three rotation domains of orthorhombic \( \kappa \)-Ga\(_2\)O\(_3\) which are not overlapping in direction with the ones of other domains. The epitaxial relationship between \( \kappa \)-Ga\(_2\)O\(_3\) with sapphire is represented as (001) \( \kappa \)-Ga\(_2\)O\(_3\) || <0001> Al\(_2\)O\(_3\), (010) \( \kappa \)-Ga\(_2\)O\(_3\) || <11–20> Al\(_2\)O\(_3\) and (100) \( \kappa \)-Ga\(_2\)O\(_3\) || <1–100> Al\(_2\)O\(_3\) [24,25].

![Fig. 1. XRD 2theta-omega scan profiles of \( \kappa \)-Ga\(_2\)O\(_3\) layers grown on (a) c-plane Sapphire (0001), (b) AlN (0001), and (c) GaN (0001). The left inset represents the RHEED images along (100) and (010) azimuthal planes where a two-dimensional film growth was witnessed. The right insets show the corresponding AFM image and schematic of the growth structure.](image1)

Fig. 2 shows the HAADF-STEM images of the cross-sectional view of (a) \( \kappa \)-Ga\(_2\)O\(_3\)/AlN interface and (b) \( \kappa \)-Ga\(_2\)O\(_3\)/GaN interface, respectively. Both images are taken along the [11–20] zone axis of AlN and GaN, as shown in the inset of the corresponding FFT (a) and (b). A crystal model of \( \kappa \)-Ga\(_2\)O\(_3\) along its [010] zone axis is overlayed on the as-grown thin film.

![Fig. 2. HAADF-STEM images of the cross-sectional view of (a) \( \kappa \)-Ga\(_2\)O\(_3\)/AlN interface and (b) \( \kappa \)-Ga\(_2\)O\(_3\)/GaN interface, respectively. Both images are taken along the [11–20] zone axis of AlN and GaN, as shown in the inset of the corresponding FFT (a) and (b). A crystal model of \( \kappa \)-Ga\(_2\)O\(_3\) along its [010] zone axis is overlayed on the as-grown thin film.](image2)
The structural properties were further cross-examined by Reflection High Energy Electron Diffraction (RHEED) and Atomic Force Microscopy (AFM). The inset of Fig. 1 represents the AFM and RHEED pattern of $\kappa$-Ga$_2$O$_3$ grown on various substrates. The bottom right inset of Fig. 1 exemplifies the corresponding growth structure schematics. A very smooth and uniform surface morphology was witnessed in the AFM images shown in top right inset for all the samples. The Root Mean Square (RMS) roughness values are 1.2 nm, 1.3 nm and 0.78 nm, respectively for the sample grown on sapphire, AlN, and GaN. The left inset represents the RHEED images along (1 0 0) and (0 1 0) azimuthal planes. A prominent streaky pattern along with 6x6 reconstruction (along (010)) were observed on $\kappa$-Ga$_2$O$_3$ grown on sapphire, AlN and GaN templates. The streaky pattern in both azimuthal images confirms the two-dimensional growth with highly single crystalline nature.

Further, the heterostructure interface quality was scrutinized by STEM. The cross-sectional STEM images of the $\kappa$-Ga$_2$O$_3$/AlN and $\kappa$-Ga$_2$O$_3$/GaN heterostructures are represented in the supplementary information S2. It exhibits a columnar growth of the material along (0 0 1) zone axis. Fig. 2 a) and 2b) shows the cross sectional high angle annular dark-field (HAADF) image of $\kappa$-Ga$_2$O$_3$/AlN and $\kappa$-Ga$_2$O$_3$/GaN heterostructure, respectively. It exemplifies a sharp interface between the heterostructure and aligned, well-ordered crystal lattice of $\kappa$-Ga$_2$O$_3$, AlN and GaN. Moreover, HR-STEM observations reveal that as-grown $\kappa$-Ga$_2$O$_3$ layer has a high crystal-quality with clear and regular atomic arrangement. The inset image represents the corresponding Fast Fourier transform (FFT) of each layer which possesses the hexagonal wurtzite features of AlN and GaN template while an orthorhombic crystal structure feature of $\kappa$-Ga$_2$O$_3$ [20]. An orthorhombic crystal model of $\kappa$-Ga$_2$O$_3$ along its [0 1 0] zone axis overlayed on the as-grown thin film is in very good match with our TEM observation, which indicates that the as-grown thin film is $\kappa$-Ga$_2$O$_3$. Therefore, their epitaxial relationships are (0 1 0) $\kappa$-Ga$_2$O$_3$ || [11 20] AlN and (0 1 0) $\kappa$-Ga$_2$O$_3$ || [11 20] GaN, respectively. No rotation domains of $\kappa$-Ga$_2$O$_3$ are observed in this experiment. Furthermore, the atomic sharp interface between $\kappa$-Ga$_2$O$_3$ and the corresponding substrate indicates that there is no transition layer during the thin film growth [23].

Band alignment of $\kappa$-Ga$_2$O$_3$/GaN and $\kappa$-Ga$_2$O$_3$/AlN heterostructures were measured by HR-XPS. The HR-XPS measurements have been performed using a Kratos Axis Ultra DLD spectrometer equipped with a monochromatic AlK$\alpha$ X-ray source (hv = 1486.6 eV) operating at 150 W, a multi-channel plate, and a delay line detector under a vacuum of ~ 10$^{-9}$ Torr. The binding energy of the C 1s peak (284.8 eV) was used as a standard reference. The valence band offset, VBO ($\Delta E_v$) of $\kappa$-Ga$_2$O$_3$ and GaN can be calculated from the following formula:

$$\Delta E_v = (E_{\text{O1s}}^{\kappa\text{-Ga}_2\text{O}_3} - E_{\text{VBM}}^{\kappa\text{-Ga}_2\text{O}_3}) - (E_{\text{N1s}}^{\text{GaN}} - E_{\text{VBM}}^{\text{GaN}}) - (E_{\text{Ga\ Auger}}^{\kappa\text{-Ga}_2\text{O}_3/\text{GaN}} - E_{\text{N1s}}^{\kappa\text{-Ga}_2\text{O}_3/\text{GaN}})$$

(1)

$E_{\text{O1s}}^{\kappa\text{-Ga}_2\text{O}_3}$ and $E_{\text{N1s}}^{\text{GaN}}$ represent the O 1s and N 1s binding energy position of bulk $\kappa$-Ga$_2$O$_3$ and GaN film, while E$_{\text{VBM}}$ denotes the valence band maximum (VBM) position of respective samples. The core level (CL) and valence band (VB) spectra of the bulk $\kappa$-Ga$_2$O$_3$ film are shown in Fig. 3 (a) and Fig. 3 (b), respectively. Fig. 3 (a) illustrates the O 1s CL spectrum acquired from the $\kappa$-Ga$_2$O$_3$ layer which manifests a single peak at 530.61 eV, corresponding to the Ga-O bond. Fig. 3(b) represents the valence band spectrum where VBM is estimated by linear extrapolating the leading edge to the baseline of the respective valence band photoelectron spectrum. The VBM of the $\kappa$-Ga$_2$O$_3$ is measured to be 2.9 eV. Thereby, the separation between the CL energies of O 1s and VBM $[\Delta E = (E_{\text{O1s}}^{\kappa\text{-Ga}_2\text{O}_3} - E_{\text{VBM}}^{\kappa\text{-Ga}_2\text{O}_3})]$ is 527.71 eV.

To determine the term in the second bracket in Eq. (1), XPS spectra of the bulk GaN were acquired. Fig. 3 (c) and 3 (d) represent the N 1s CL and VB spectra of GaN. The N 1s spectrum is deconvoluted into three peaks to exactly calculate the binding energy position of the Ga-N bond. The first two peaks positioned at 383.4 eV and 395.55 eV correspond to Ga-Auger peaks. However, the deconvoluted third peak corresponds to the Ga-N
bond from the N1s spectrum which is positioned at 397.15 eV. Therefore, the difference between the CL energies of N1s and VBM \(\Delta E = (E_{\text{N1s}} - E_{\text{VBM}})\) is determined to be 394.55 eV. The details of the fitting parameters and binding energy positions are tabulated in Table 1. The last bracket term in Eq. (1) represents the CL separation between the O1s and N1s peaks from the XPS measurement of \(\kappa\)-Ga2O3/GaN heterostructure. Fig. 3(e) and 3(f) show the O1s and N1s CLs, which is originated from \(\kappa\)-Ga2O3/GaN heterostructure, respectively. Fig. 3(e) represents the single peak of O1s positioned at 530.81 eV. Further, the N1s CL spectrum of the heterostructure was deconvoluted to three peaks, and the peak corresponding to the Ga-N bond from the N1s spectrum is positioned at 397.43 eV. Thereby, the energy difference between the CLs \(E_{\text{Ga-N}}\) is 133.38 eV. Thus, the VBO is calculated to be \(-0.22 \pm 0.07\) eV (the error value is calculated by considering the fitting error and discrepancy in the repeated measurements). The conduction band offset, CBO (\(\Delta E_c\)) for \(\kappa\)-Ga2O3/GaN heterostructure is determined by substituting the VBO and electronic bandgap \((E_g)\) values of \(\kappa\)-Ga2O3 and GaN in equation (2),

\[
\Delta E_c = E_{\text{Ga-N}} - E_{\text{O}}^{\text{GaN}} + \Delta E_v \tag{2}
\]

\(E_{\text{Ga-N}}\) and \(E_{\text{O}}^{\text{GaN}}\) were measured from the transmittance tauc plot, and it is calculated to be 5.0 eV and 3.44 eV, respectively (supplementary information S3). Hence, the CBO is estimated to be 1.38 ± 0.07 eV. These experimentally determined parameters are assimilated into a band alignment diagram as represented in Fig. 5, which divulges that the \(\kappa\)-Ga2O3/GaN heterostructure forms a type-I band alignment. This peculiar band alignment with high CBO and huge polarization difference of the materials could form the two dimensional electron gas (2DEG) at the interface. The dependency of polarity (Ga and N-polar) on the band alignment of \(\kappa\)-Ga2O3/GaN will be an interesting study to be carried out in the future. Further, the band alignment of \(\kappa\)-Ga2O3 with AlN WBG semiconductor is depicted below.

The valence band offset of \(\kappa\)-Ga2O3 with AlN can be calculated from the following formula:

\[
\Delta E_v = E_g^{\text{AIN}} - E_g^{\text{Ga2O3}} - (E_{\text{VBM}}^{\text{AIN}} - E_{\text{VBM}}^{\text{Ga2O3}}) \tag{3}
\]

\(E_{\text{VBM}}^{\text{Ga2O3}}\) and \(E_{\text{VBM}}^{\text{AlN}}\) are the valence band maximums of \(\kappa\)-Ga2O3 and AlN, respectively.

### Table 1

<table>
<thead>
<tr>
<th>Sample</th>
<th>State</th>
<th>Binding energy (eV)</th>
<th>Bonding</th>
<th>FWHM (eV)</th>
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<td>530.61</td>
<td>Ga-O</td>
<td>1.7</td>
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<tr>
<td></td>
<td>VBM</td>
<td>2.9</td>
<td></td>
<td></td>
</tr>
<tr>
<td>GaN</td>
<td>N1s</td>
<td>397.15</td>
<td>Ga-N</td>
<td>1.2</td>
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<tr>
<td></td>
<td>VBM</td>
<td>2.6</td>
<td></td>
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<tr>
<td>(\kappa)-Ga2O3/GaN</td>
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<td>Ga-O</td>
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<tr>
<td></td>
<td>N1s</td>
<td>397.43</td>
<td>Ga-N</td>
<td>1.2</td>
</tr>
</tbody>
</table>

Fig. 3. XPS core level spectra of (a) Ga3d and (b) valence band spectra of the \(\kappa\)-Ga2O3; core level spectra of (c) Al2p and (d) valence band spectra of the AlN; core level spectra of (e) Ga3d and (f) Al2p of the \(\kappa\)-Ga2O3/AlN heterostructure.

![Fig. 5. The band alignment diagram of the \(\kappa\)-Ga2O3/GaN heterojunction, along with that of \(\kappa\)-Ga2O3/AlN.](image-url)
ΔE = \left( E_{Ga2O3}^{VBM} - E_{Al2p}^{VBM} \right) - \left( E_{Ga2O3}^{CBM} - E_{Al2p}^{CBM} \right)

(3)

E_{Ga3d} and E_{Al2p} represent the peak binding energy position of the Ga3d and Al2p CL spectra, respectively, while E_{VBM} denotes the valence band maximum position. Fig. 4 (a) and (b) represent the Ga3d CL and valence band spectra of κ-Ga2O3 bulk film. In Fig. 4 (a), Ga3d binding energy is positioned at 20.15 eV with a Gaussian peak FWHM value of 1.3 eV. The valence band maximum value was evaluated to be 3.02 eV from the valence band spectra (Fig. 4 (b)). From Fig. 4 (c) and (d) (for bulk AlN film), Al2p position and VBM values were calculated to be 73.1 eV and 2.4 eV, respectively. These values are in good agreement with our previous report [21]. Fig. 4(e) and 4(f) show the Ga3d and Al2p CLs, which originated from heterostructure, respectively. The details of the fitting parameters and binding energy positions are tabulated in Table 2. Putting all these experimental values in equation (3), the VBO for κ-Ga2O3/AlN heterostructure is determined to be 0.08 ± 0.07 eV.

The conduction band offset (ΔE\textsubscript{C}) for κ-Ga2O3/AlN heterostructure is determined by substituting the VBO and the electronic bandgap values in the equation, ΔE\textsubscript{C} = E_{Ga2O3}^{CBM} - E_{AlN}^{CBM} + ΔE\textsubscript{V}. Thereby, CBO is calculated to be 1.04 ± 0.07 eV, where the E_{AlN} is measured to be 6.12 eV from the transmittance tauc plot (supplementary information S3).

Fig. 5 represents the band alignment of both κ-Ga2O3/AlN and κ-Ga2O3/AlN heterostructures, which forms a type-I band alignment. However, in our previous study, the monolinic β-Ga2O3/AlN heterostructure forms a type-II band alignment with VBO of ~0.55 ± 0.05 eV [21]. While the monolinic β-Ga2O3/AlN heterostructure forms a type-I band alignment with a high VBO of 1.4 ± 0.08 eV [27]. This particular difference in electronic band alignment property of the same material with different phases could be utilized for various applications. In the present study, the formation of type-I is interesting for developing high electron mobility transistors for two reasons. First of all, κ-Ga2O3 possesses spontaneous polarization and ferroelectric properties so that the polarization can be tuned by applying bias voltage [19]; however, other phases of Ga2O3 do not show this unique material property. Secondly, integrating well-established nitride polarization materials with ferroelectric κ-Ga2O3 will allow more space for tunability of the polarization difference between the materials and thereby carrier concentration within the interface. There are a few simulation studies which explain the huge interface charges at the κ-Ga2O3/(Al, GaN) interface [15,17]. Moreover, the experimental results of the band alignment study suggest that the electron affinity of κ-Ga2O3 will be lesser than the electron affinity of β-Ga2O3 (≤ 3.15 eV).

4. Conclusions

In conclusion, orthorhombic Ga2O3 was deposited on GaN and AlN templates to form a heterojunction. XPS measurements were carried out to determine the VBO and CBO of κ-Ga2O3/GaN and κ-Ga2O3/AlN heterostructures. κ-Ga2O3 forms a type-I heterojunction with GaN and AlN with high CBO values of 1.38 ± 0.07 eV and 1.04 ± 0.07 eV, respectively. The existence of spontaneous polarization property along with unique band alignment suggests the possibility of 2DG at the heterointerface. Therefore, this work could lead to the development of efficient power devices with tunable carrier concentrations.

CRediT authorship contribution statement

Shibin Krishna: Conceptualization, Investigation, Methodology, Data curation, Formal analysis, Writing – original draft, Writing – review & editing. Yi Lu: Methodology. Che-Hao Liao: Methodology. Vishal Khandelwal: Xiaohang Li: Conceptualization, Investigation, Funding acquisition, Project administration, Resources, Supervision, Validation, Writing – review & editing.

Declaration of Competing Interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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Appendix A. Supplementary material

Supplementary data to this article can be found online at https://doi.org/10.1016/j.apsusc.2022.153901.

References


